Bad Cavities for Good Memories: Storing Broadband Photons with Low Noise

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INTRODUCTION

Linear optical logic elements are non-deterministic, and cannot be efficiently concatenated unless a multiplexing strategy is employed to actively pick out successful gate operations. Quantum memories promise scalability via temporal multiplexing, where successful outputs are stored until all parts of a photonic circuit have executed [1].

Off-resonant Raman memory protocols, in which a strong *control* field mediates the two-photon absorption of a signal in a Λ -type ensemble, enable multimode spectral/temporal control of long pulses [2, 3], or extremely broadband storage of short pulses [4, 5].

But the strong control field can drive spontaneous Raman scattering that produces spurious excitations in the memory, which are then retrieved as noise [6, 7]. This *fourwave mixing* (FWM) noise degrades the non-classical character of single-photon states stored in the memory [8, 9], and is the last remaining roadblock to developing efficient, configurable, broadband quantum photonic memories capable of operating at room temperature or in the solid state.

PROPOSAL

Here, we propose a method to suppress FWM noise by means of a low-finesse optical cavity. In a non-degenerate Λ -system, spontaneous Raman emission from the ground state is Stokes-shifted in frequency, being distinguishable from both the control field and the signal to be stored. Placing the memory inside a cavity tuned so that the on-axis Stokes emission is anti-resonant with the cavity mode provides a suppression of the FWM noise proportional to the cavity finesse \mathcal{F} . While $\mathcal{F} > 10$ is generally sufficient to remove the influence of the FWM noise, the protocol is compatible with a broad acceptance bandwidth for the memory.

The scheme has several other advantages: first, the effective length of the memory is enhanced by the cavity, since the signal field makes multiple passes through the memory medium [10]. This, combined with the cavity-enhanced control field, confined to a smaller mode volume than in free space, enables a higher memory efficiency in a smaller footprint. Interestingly, the cavity also constrains the memory interaction to a single longitudinal spatial mode, which in turn means that only a single optical mode, of arbitrarily configurable shape, is coupled to the memory. This has applications in quantum optical signal processing [11].

EXPERIMENT

We also describe an experimental implementation of the above scheme, in which a small caesium vapour-cell is mounted inside a birefringent ring cavity, stabilised via two-colour Hänsch-Couillaud phase locking, with a resonance blue-detuned from the Cs D_2 line by 15 GHz. We demonstrate, for the first time, cavity-enhanced Raman storage and retrieval of sub-ns coherent states. Finally we present the noise characteristics of the cavity memory, and discuss routes towards an optimised design.

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- [1] J. Nunn, N. Langford, W. Kolthammer, T. Champion, M. Sprague, P. Michelberger, X.-M. Jin, D. England, and I. Walmsley, Phys. Rev. Lett. **110**, 133,601 (2013).
- [2] M. Hosseini, B. M. Sparkes, G. Hetet, J. J. Longdell, P. K. Lam, and B. C. Buchler, Nature 461, 241–245 (2009).
- [3] M. Hosseini, G. Campbell, B. Sparkes, P. Lam, and B. Buchler, Nature Physics 7, 795–799 (2011).
- [4] K. Reim, P. Michelberger, K. Lee, J. Nunn, N. Langford, and I. Walmsley, Phys. Rev. Lett. 107, 053,603 (2011).
- [5] D. England, P. Bustard, J. Nunn, R. Lausten, and B. Sussman, Phys. Rev. Lett. **111**, 243,601 (2013).
- [6] D. G. England, K. A. Fisher, J.-P. W. MacLean, P. J. Bustard, R. Lausten, K. J. Resch, and B. J. Sussman, arXiv:1409.2892 (2014).
- [7] P. Michelberger, T. Champion, M. Sprague, K. Kaczmarek, M. Barbieri, X. Jin, D. England, W. Kolthammer, D. Saunders, J. Nunn *et al.*, arXiv:1405.1470 (2014).
- [8] N. Lauk, C. O'Brien, and M. Fleischhauer, Phys. Rev. A 88, 013,823 (2013).
- [9] M. Dabrowski, R. Chrapkiewicz, and W. Wasilewski, Op. Ex. 22, 26,076–26,091 (2014).
- [10] A. V. Gorshkov, A. André, M. D. Lukin, and A. S. Sørensen, Phys. Rev. A 76, 033,804 (2007).
- [11] D. V. Reddy, M. G. Raymer, C. J. McKinstrie, L. Mejling, and K. Rottwitt, Op. Ex. 21, 13,840–13,863 (2013).